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**“Relative h-principles” for negative scalar curvature.**

Are Lorentzian metrics of negative scalar curvature ubiquitous?

... a priori only loosely related, but there is a common underlying theme:

nonintegrable distributions. (*distribution on  $M$  = sub vector bundle of  $TM$* )

## §1. The dominant energy condition

... is an inequality for the Ricci curvature of a Lorentzian  $n$ -manifold.

Sign convention: *Lorentzian* means  $-+++ \dots +$ , i.e. index 1, signature  $(n-1, 1)$ .

$v \in TM$  is *spacelike*  $\Leftrightarrow g(v, v) > 0$

$v \in TM$  is *timelike*  $\Leftrightarrow g(v, v) < 0$

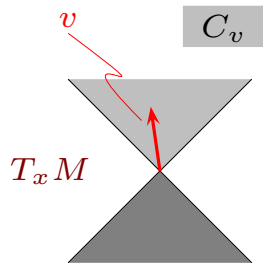
$v \in TM$  is *lightlike*  $\Leftrightarrow g(v, v) = 0$  and  $v \neq 0$ .

### **Definition.**

Let  $(M, g)$  be a Lorentzian manifold, let  $\Lambda \in \mathbb{R}$ .

Consider the *energy-momentum tensor*  $T := \text{Ric}_g - \frac{1}{2}\text{scal}_g g + \Lambda g$  with respect to the “cosmological constant”  $\Lambda$ .

For a timelike or lightlike  $v \in T_x M$ , let  $C_v$  denote the closure of the **connected component** of  $\{w \in T_x M \mid w \text{ is timelike or lightlike}\}$  which contains  $v$ .



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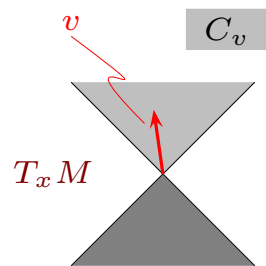
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$(M, g)$  satisfies *the dominant energy condition w.r.t.  $\Lambda$*

iff  $-T_a^b v^a \in C_v$  for every timelike (or lightlike)  $v \in TM$ .

$(M, g)$  satisfies *the strict dominant energy condition w.r.t.  $\Lambda$*

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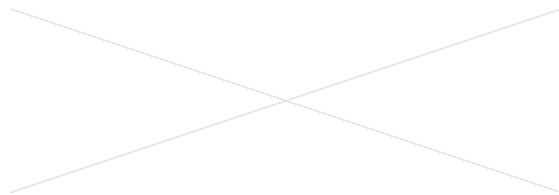


Does every manifold which admits a Lorentzian metric at all admit one which satisfies the dominant energy condition?

The answer is *yes* in most cases:

**Theorem 1.** Let  $(M, g)$  be a connected Lorentzian manifold of dimension  $n \geq 4$ , let  $\Lambda \in \mathbb{R}$ , let  $K$  be a compact subset of  $M$ .

If  $n = 4$ , assume that  $(M, g)$  is time- and space-orientable, and that either  $\partial M \neq \emptyset$ , or  $M$  is closed with intersection form signature divisible by 4.



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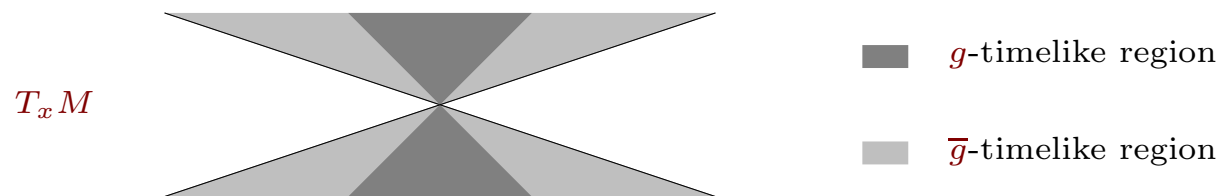
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Then there exists a Lorentzian metric  $\bar{g}$  on  $M$  such that

- $\bar{g}$  satisfies the *strict dominant energy condition with respect to  $\Lambda$  on  $K$* ;
- *every  $g$ -timelike or  $g$ -lightlike vector in  $TM$  is  $\bar{g}$ -timelike.*

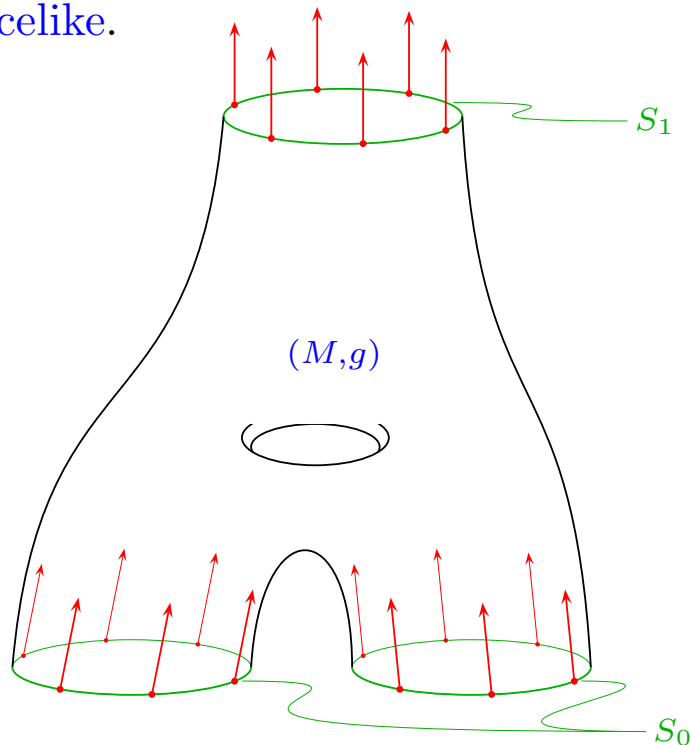


A classical topic in General Relativity is the problem of *topology change*:

**Definition.** Let  $S_0, S_1$  be  $(n - 1)$ -dimensional closed manifolds.

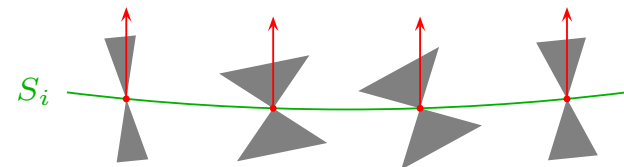
A *weak Lorentz cobordism* between  $S_0$  and  $S_1$  is a compact Lorentzian  $n$ -manifold  $(M, g)$  whose boundary is the disjoint union  $S_0 \sqcup S_1$ , such that  $M$  admits a timelike vector field which is inward-directed on  $S_0$  and outward-directed on  $S_1$ .

A *Lorentz cobordism* is a weak Lorentz cobordism  $(M, g)$  such that  $\partial M$  is  $g$ -spacelike.

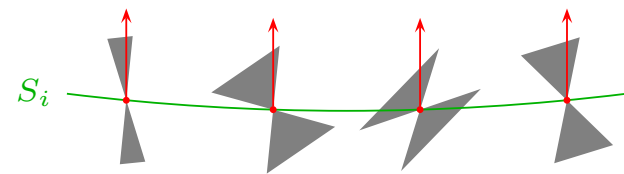


↑↑↑ timelike vector field on  $(M, g)$

lightcones at the boundary:



Lorentz cobordism



weak Lorentz cobordism

**Fact.**  $S_0, S_1$  are weakly Lorentz cobordant **iff** they are Lorentz cobordant.

**Theorem [Tipler 1977].** *If there exists a Lorentz cobordism  $(M, g)$  between closed 3-manifolds  $S_0$  and  $S_1$  such that  $\text{Ric}_g(v, v) > 0$  holds for all lightlike  $v \in TM$ , then  $S_0$  and  $S_1$  are diffeomorphic.*

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**Theorem 2.** Let  $S_0, S_1$  be arbitrary closed orientable 3-manifolds, let  $\Lambda \in \mathbb{R}$ . Then there exists a weak Lorentz cobordism  $(M, g)$  between  $S_0$  and  $S_1$  which satisfies the strict dominant energy condition with respect to  $\Lambda$ ; in particular,  $\text{Ric}_g(v, v) > 0$  holds for all lightlike  $v \in TM$ .

## §2. Prescribed scalar curvature

- Well-known:
- Every manifold of  $\text{dim.} \geq 3$  admits a Riemannian  $g$  with  $\text{scal}_g < 0$ .
  - Some closed manifolds, e.g. the  $n$ -dimensional torus, do **not** admit a Riemannian metric  $g$  with  $\text{scal}_g > 0$ .

Which smooth functions are scalar curvatures of Lorentzian metrics?

And what about pseudo-Riemannian metrics of index  $\geq 2$ ?

**Convention.** In the rest of this §2, we assume that

*all manifolds are connected and **compact**, possibly with boundary.*

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**Theorem 3.** Let  $(M, g)$  be a semi-Riemannian manifold of signature  $(p, q)$ , where  $p \geq 3$  and  $q \geq 1$ . Let  $s \in C^\infty(M, \mathbb{R})$ .

*Then  $M$  admits a metric  $\bar{g}$  of index  $q$  and scalar curvature  $s$  such that every  $g$ -timelike or  $g$ -lightlike vector in  $TM$  is  $\bar{g}$ -timelike.*

(No restriction on  $s$ ! Take e.g.  $(M, g)$  to be a flat Lorentzian  $n$ -torus and  $s \equiv 1$ .)

Let  $\mathcal{M}_q(M)$  denote the set of semi-Riemannian metrics of index  $q$  on a manifold  $M$ . We equip  $\mathcal{M}_q(M)$  (viewed as a set of sections in the bundle  $\text{Sym}^2 T^*M \rightarrow M$ ) with the compact-open topology.

**Theorem 4.** Let  $M$  be an  $n$ -manifold with  $n \geq 3$ , or  $n = 2$  and  $\partial M \neq \emptyset$ .

Let  $1 \leq q \leq n - 1$ . Then for every  $s \in C^\infty(M, \mathbb{R})$ ,

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When  $n = 2$  and  $\partial M = \emptyset$ , an obstruction occurs:

The Gauß/Bonnet formula for closed Lorentzian surfaces [Avez 1962] says that

$$\int_M \text{scal}_g \mu_g = 0 .$$

So if  $s = \text{scal}_g$  for some metric  $g$ , then  $s$  is identically 0 or changes its sign.

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**Theorem 5.** Let  $M$  be the 2-dimensional torus or the Klein bottle

(these are the only closed nonempty 2-manifolds which admit a Lorentzian metric).

A function  $s \in C^\infty(M, \mathbb{R})$  is the scalar curvature of a Lorentzian metric on  $M$  if and only if  $s$  is either identically 0 or changes its sign.

### §3. Relative h-principles for $\text{scal} < 0$

(The convention from §2 is no longer valid: manifolds are not necessarily compact.)

#### **Definition of the fine $C^{k,\alpha}$ -topology.**

Let  $k \in \mathbb{N}$  and  $\alpha \in [0, 1]$ .

We define the *fine  $C^{k,\alpha}$ -topology* on the set of  $C^\infty$  sections in  $\text{Sym}^2 T^*M \rightarrow M$ :

Let  $h$  be a Riem. metric on  $M$ , let  $K \subseteq (M, h)$  be compact & geodesically convex.

For  $x, y \in K$ , let  $\tau_{x,y} :=$  parallel transport along the unique geodesic  $x \rightarrow y$  in  $K$ .

We have the Hölder  $C^{k,\alpha}$ -norm  $\|\cdot\|_{C^{k,\alpha}(K;h)}$  for  $C^\infty$  tensor fields  $\eta, \sigma$  on  $K$ :

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$$D^j \eta := \underbrace{\nabla^h \dots \nabla^h}_{j \text{ derivatives}} \eta \quad \|\sigma\|_{C^0(K;h)} := \max\{|\sigma(x)|_h \mid x \in K\}$$

$$[\sigma]_{C^{0,\alpha}(K;h)} := \sup \left\{ \frac{|\tau_{x,y}(\sigma(x)) - \sigma(y)|_h}{\text{dist}_h(x,y)^\alpha} \mid x, y \in K, x \neq y \right\}$$

$$\|\eta\|_{C^{k,\alpha}(K;h)} := \sum_{j=1}^k \|D^j \eta\|_{C^0(K;h)} + [D^k \eta]_{C^{0,\alpha}(K;h)}$$

The **topology** induced by  $\|\cdot\|_{C^{k,\alpha}(K;h)}$  **does not depend on  $h$ .**

The *fine*  $C^{k,\alpha}$ -topology on  $\mathcal{M}_q(M)$  is defined via neighbourhood bases  $\mathfrak{B}_g$ :

We choose • a Riemannian metric  $h$  on  $M$ ;

- a locally finite family  $(K_i)_{i \in \mathbb{N}}$

of compact & geodesically convex sets  $\subseteq M$  whose interiors cover  $M$ .

For every sequence  $\varepsilon = (\varepsilon_i)_{i \in \mathbb{N}}$  in  $\mathbb{R}_{>0}$  and every  $g \in \mathcal{M}_q(M)$ , we define

$$\mathcal{N}_{g,\varepsilon} := \{ \bar{g} \in \mathcal{M}_q(M) \mid \forall i \in \mathbb{N}: \|\bar{g} - g\|_{C^{k,\alpha}(K_i;h)} < \varepsilon_i \}$$

$$\mathfrak{B}_g := \{ \mathcal{N}_{g,\varepsilon} \mid \varepsilon = (\varepsilon_i)_{i \in \mathbb{N}}, \varepsilon_i \in \mathbb{R}_{>0} \} .$$

The topology defined by the nbhd. bases  $\mathfrak{B}_g$  does not depend on  $h$  or  $(K_i)_{i \in \mathbb{N}}$ .

In the Riemannian case  $q = 0$ , J. Lohkamp proved the following:

**Theorem [Lohkamp 1995: the “relative  $C^0$ -dense h-principle for  $\text{scal} < c$ ”].**

Let  $(M, g)$  be a Riemannian manifold of dimension  $p \geq 3$ .

Let  $A$  be a closed subset of  $M$ , let  $c \in \mathbb{R}$  be such that  $\text{scal}_g < c$  holds on  $A$ .

Let  $\mathcal{U}$  be a neighbourhood of  $g$  in  $\mathcal{M}_0(M)$  with respect to the fine  $C^0$ -topology.

*Then there exists a metric  $\bar{g} \in \mathcal{U}$  with  $\text{scal}_{\bar{g}} < c$  such that  $\bar{g} = g$  holds on  $A$ .*

In particular (take  $A = \emptyset$ ,  $c = 0$ ), the set of negative scalar curvature metrics is dense in the space of Riemannian metrics with respect to the fine  $C^0$ -topology.

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(Lohkamp also proved an analogous theorem with  $\text{Ric}$  instead of  $\text{scal}$ .)

Lohkamp showed that **the theorem fails with  $C^1$  or  $C^{0,1}$  instead of  $C^0$ .**

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**Theorem 6 [the semi-Riem. relative  $C^{0,\alpha}$ -dense h-principle for  $\text{scal} < c$ ].**

Let  $(M, g)$  be a **semi-Riemannian manifold of signature  $(p, q)$** , where  $p \geq 3$ .

Let  $0 \leq \alpha < 1$ . Let  $A$  be a closed subset of  $M$ , let  $c \in \mathbb{R}$  with  $\text{scal}_g < c$  on  $A$ .

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## §4. The common theme of the proofs:

### “stretching” metrics along nonintegrable distributions

Techniques in the proofs that we will *not* discuss here:

- prescribed scalar curvature  $\rightsquigarrow$  differential *equation* (underdetermined);  
trick: restrict underdetermined PDE to *elliptic* PDE
- Gromov’s h-principle for ample partial differential relations  
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Focus on *inequalities*:

- How to prove existence of positive scalar curvature metrics?  
Why does the proof fail in the Riemannian case?
- How to prove ubiquity of negative scalar curv. metrics (relative h-principle)?
- How to construct dominant energy metrics?

The basic construction is the same in each case:

*stretching along a nonintegrable distribution.*

Recall that a  $k$ -plane distribution  $H$  on a manifold  $M$  (i.e. a sub vector bundle of rank  $k$  of  $TM$ ) is *integrable* iff one of the following equivalent statements holds:

1. There exists a  $k$ -dimensional foliation  $F$  of  $M$  such that  $\forall x \in M: H_x = T_x F$ .
2.  $\forall v, w \in \Gamma(H) : [v, w] \in \Gamma(H)$ .

We're interested in distributions  $H$  which are *nonintegrable* in the following sense:

$$\forall x \in M : \exists v, w \in \Gamma(H) : [v, w](x) \notin H_x .$$

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Let  $(M, g)$  be a semi-Riemannian manifold, let  $f \in C^\infty(M, \mathbb{R}_{>0})$ , let  $TM = V \oplus H$  be a  $g$ -orthogonal splitting. We “stretch”  $g = g_V \oplus g_H$  along  $V$ :

$$\bar{g} = \left( \frac{1}{f^2} g_V \right) \oplus g_H \ .$$

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$$\bar{g} = \left( \frac{1}{f^2} g_V \right) \oplus g_H .$$

The scalar curvature of  $\bar{g}$  looks like this (with  $q = \text{index}(g)$ ):

$$\begin{aligned} \text{scal}_{\bar{g}} = & \text{scal}_g + \frac{2q}{f} \Delta_g^H(f) + 2(q-1)f \Delta_g^V(f) - \frac{q(q+3)}{f^2} |df|_{g,H}^2 - q(q-1) |df|_{g,V}^2 \\ & + 2(q-2)f \langle \text{div}_g^H, df \rangle_{g,V} + \frac{2(q+1)}{f} \langle \text{div}_g^V, df \rangle_{g,H} \\ & + (1-f^2) \xi_g^V - \frac{f^4 - f^2}{4} |\mathcal{T}w_g^V|_g^2 - \frac{1-f^2}{4f^2} |\mathcal{T}w_g^H|_g^2 . \end{aligned}$$

## Dominant energy: Idea of proof.

1. Choose a  $g$ -orthogonal splitting  $TM = V \oplus H$  with  $V$  timelike,  $H$  spacelike. Show that this can be done with **nonintegrable**  $H$ .
2. Using a very **small constant**  $f > 0$ , stretch  $g$  along  $V \rightsquigarrow$  d.e.c. holds. If  $f < 1$ , then the lightcones of the new metric are wider than those of  $g$ .

### Dominant energy: Idea of proof.

1. Choose a  $g$ -orthogonal splitting  $TM = V \oplus H$  with  $V$  timelike,  $H$  spacelike. Show that this can be done with **nonintegrable**  $H$ .
2. Using a very **small constant**  $f > 0$ , stretch  $g$  along  $V \rightsquigarrow$  d.e.c. holds. If  $f < 1$ , then the lightcones of the new metric are wider than those of  $g$ .

### Relative $C^{0,\alpha}$ -dense h-principle: Idea of proof.

1. Choose a  $g$ -orthogonal splitting  $TM = V \oplus H$  with  $V$  timelike,  $H$  spacelike. Keep the metric fixed on  $V$ , modify it on  $H$  as follows.
2. Cover  $M$  by compact balls. On each ball  $B_i$ , choose a 1-parameter family of  $g$ -orth. splittings  $H = V'_c \oplus H'_c$  s.t.  $H'_c$  becomes “**more nonintegrable**” as  $c \rightarrow \infty$ .
3. Choose cutoff  $\beta \in C^\infty(M, [0, 1])$  which is 0 outside  $B_i$  and is 1 on a slightly smaller ball  $B'_i$ . Using  $f = 1 - a\beta \in C^\infty(M, \mathbb{R}_{>0})$  with a **small constant**  $a > 0$ , stretch  $g$  along  $V'_c$ .

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4. Do this with  $c = c(a)$ , s.t.  $\lim_{a \rightarrow 0} c(a) = \infty$ , but  $c(\cdot)$  **not faster increasing than necessary**. For  $a \rightarrow 0$ , the new metric is  $C^{0,\alpha}$ -close to  $g$ , but has  $\text{scal} \ll 0$  on  $B'_i$ .
5. Do this for every  $B_i$ .

## Prescribed scalar curvature: Idea of proof (very roughly).

1. Solve the elliptic equation (mentioned above, but not explained) via the **method of sub- and supersolutions**. The equation depends on data  $g$  and  $H$ . Construct  $g$  and  $H$  such that sub- and supersolutions exist, as follows.
2. Choose a splitting  $TM = V \oplus H$  with  $V$  timelike,  $H$  spacelike. Show that this can be done with  $H$  being **nonintegrable** (almost everywhere).
3. **Change the metric  $g$  on  $H$  as in the h-principle proof, creating so much negative curvature** that (say) the constant **0.1 becomes a supersolution**.
4. Because  $H$  is nonintegrable, sufficiently **small constants  $f < 1$  are subsolutions**. (If  $H$  is nonintegrable only *almost* everywhere, deal with the problem set by making  $f$  slightly nonconstant near it.)

## The elliptic equation

$$\begin{aligned}
 0 = & +2\Delta_g(f) + a_{n,q}(f)|df|_g^2 + b_{n,q}(f)|df|_{g,V}^2 \\
 & + \frac{2(1+f^2)}{f^2} \langle \operatorname{div}_g^V, df \rangle_{g,H} + 2(1+f^2) \langle \operatorname{div}_g^H, df \rangle_{g,V} \\
 & + \frac{(1+f^2)^2}{2f^3} |\operatorname{Twist}_H|_g^2 - \frac{f(1+f^2)^2}{2} |\operatorname{Twist}_V|_g^2 + \frac{(1+f^2)^2}{f} \xi_{g,V} \\
 & + \frac{1+f^2}{f} \operatorname{scal}_g - f^{\frac{2q}{n-1}-1} (1+f^2)^{1-\frac{1}{n-1}} s .
 \end{aligned}$$

The functions  $a_{n,q}, b_{n,q} \in C^\infty(\mathbb{R}_{>0}, \mathbb{R})$  are defined by

$$\begin{aligned}
 a_{n,q}(x) & := \frac{((q-1)^2 - (n-1)(q+3))x^4 - 2(q-1)(n-1-q)x^2 - q(n-1-q)}{(n-1)x^3(1+x^2)} , \\
 b_{n,q}(x) & := \frac{(q-1)(n-q)x^4 + 2(q-1)(n-1-q)x^2 + q(n-1-q)}{(n-1)x^3} .
 \end{aligned}$$

The first-order terms are essentially irrelevant.